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High photosensitivity and broad spectral response of multi-layered germanium sulfide transistors†

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In this paper, we report the optoelectronic properties of multi-layered GeS nanosheet (~28 nm thick)-based field-effect transistors (called GeS-FETs). The multi-layered GeS-FETs exhibit remarkably high photoresponsivity of $R_\lambda \sim 206 \text{ A W}^{-1}$ under $1.5 \mu\text{W cm}^{-2}$ illumination at $\lambda = 633 \text{ nm}$, $V_g = 0 \text{ V}$, and $V_{ds} = 10 \text{ V}$. The obtained $R_\lambda \sim 206 \text{ A W}^{-1}$ is excellent as compared with a GeS nanoribbon-based and the other family members of group IV–VI-based photodetectors in the layered-materials realm, such as GeSe and SnS₂. The gate-dependent photoresponsivity of GeS-FETs was further measured to be able to reach $R_\lambda \sim 655 \text{ A W}^{-1}$ operated at $V_g = -80 \text{ V}$. Moreover, the multi-layered GeS photodetector holds high external quantum efficiency (EQE $\sim 4.0 \times 10^4\%$) and specific detectivity ($D^* \sim 2.35 \times 10^{13}$ Jones). The measured D^* is comparable to those of the advanced commercial Si- and InGaAs-based photodiodes. The GeS photodetector also shows an excellent long-term photoswitching stability over a long period of operation (>1 h). These extraordinary properties of high photocurrent generation, broad spectral range, and long-term stability make the GeS-FET photodetector a highly qualified candidate for future optoelectronic applications.

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Introduction

Layered materials are attractive building blocks in the modern electronic world. The exploration of such novel layered structures is especially interesting in optoelectronics for applications such as in transistors, photodetectors, ultrafast lasers, touch panels, and optical modulators to achieve the ever increasing demands for high-speed performance, less-power consumption, and easy-to-fabricate flexible devices.^{1–4} Compared with zero- and one-dimensional materials, by virtue of their excellent electron mobility and large surface area, two-dimensional (2D) materials are well suited for easy integration

into the complex electronic circuitry that pervades the modern electronics industry.^{5,6}

The discovery of graphene has ignited tremendous development in materials science because of its unique, novel, and exciting electrical and mechanical properties. However, the absence of a band gap and weak light absorption in graphene limit its advantageous optoelectronic applications, therefore paving the way for other 2D layered semiconducting materials with versatile electronic and optical applications. To date, various layered materials, such as transition metal dichalcogenides (*e.g.*, MoS₂,⁷ MoSe₂,⁸ WS₂,⁹ and MoTe₂¹⁰) and other multi-layered 2D crystals (*e.g.*, GaS,¹¹ GaSe,¹² and InSe¹³) have attracted significant attention and have been investigated. In particular, the layered chalcogenides of group IV–VI (GeS,¹⁴ GeSe,¹⁵ SnS,¹⁶ SnSe,¹⁷ *etc.*) have attracted wide research interest because of their suitable band gaps, high absorption coefficients, anisotropic optical/electrical properties, high thermal stability, low toxicity, and environmental sustainability. For example, single-crystal germanium selenide (GeSe) nanosheet-based photodetectors were demonstrated to possess excellent photoresponsivity ($R_\lambda \sim 3.5 \text{ A W}^{-1}$) and photoconductive gain ($\sim 5.3 \times 10^2\%$).¹⁸ These encouraging results motivated us to investigate the photoresponsive properties of germanium sulfide (GeS), which have not been explored elaborately so far, despite the band gap of GeS lying in the convenient visible

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region and the strong absorption of GeS spanning in a wide spectral range.

GeS is a p-type semiconductor of an orthorhombic structure with several direct and indirect transitions very close in energy competing for the band gap at ~ 1.55 – 1.65 eV, rendering GeS a strong absorber from the visible to near-infrared spectral regions.^{19–22} Herein, we report the electrical and optoelectronic characterization of GeS nanosheets, which were exfoliated mechanically from a bulk GeS crystal. Multi-layered GeS nanosheets (~ 28 nm in thickness) were then used as a conducting channel to fabricate field-effect transistors (FETs). The multi-layered GeS-FETs exhibit remarkably high $R_{\lambda} \sim 206$ A W^{-1} under illumination of 1.5 $\mu\text{W cm}^{-2}$ at $\lambda = 633$ nm ($V_g = 0$ V, $V_{ds} = 10$ V) and possess robust optoelectronic merits to serve as photodetectors of high photocurrent generation. The obtained $R_{\lambda} \sim 206$ A W^{-1} is excellent in the photocurrent generation as compared with a GeS nanoribbon-based photodetector and the other family members of other group IV–VI photodetectors, such as GeSe and SnS₂ (see ESI Table S1†). Moreover, the multi-layered GeS photodetector exhibits high external quantum efficiency (EQE $\sim 4.0 \times 10^4\%$) and specific detectivity ($D^* \sim 2.35 \times 10^{13}$ Jones). The measured D^* is comparable to the advanced few-layered GaS ($\sim 10^{13}$ Jones)¹¹ and commercial Si (10^{13} Jones)²³ based photodiodes. Furthermore, the D^* value can be increased by simply tuning the external gate voltage. Light-induced current *vs.* elapsed time (I_{ds} - t) measurements of the multi-layered GeS-FETs reveal the excellent photoresponse time and photoswitching stability of the materials over a long period of operation. These figures-of-

merit demonstrate that GeS nanosheets hold great potential for optoelectronic applications, especially for fabricating photodetectors with high photocurrent generation, fast response, and long-term stability.

Results and discussion

The bulk GeS crystals used in this study were grown *via* a chemical vapor transport (CVT) method²⁴ with the synthetic details being described in the Experimental section. Fig. 1a shows the three-dimensional (3D) crystal structure of GeS, which emphasizes the stacking order and van der Waals interaction among individual sheets in the GeS crystal. The X-ray diffraction (XRD) pattern of the bulk GeS crystal in Fig. 1b confirms that the GeS crystal belongs to the orthorhombic *Pbnm* space group (JCPDS no. 00-051-1168) with lattice constants of $a = 4.30$ Å, $b = 10.47$ Å, and $c = 3.64$ Å. The inset of Fig. 1b presents the selected area electron diffraction (SAED) pattern which exhibits the single-crystallinity of the GeS nanosheets oriented along the [100] direction.²⁵ In addition, the binding energies of Ge (3d) at 29.6 eV, S (2p_{3/2}) at 160.2 eV, and S (2p_{1/2}) at 161.4 eV, observed by X-ray photoelectron spectroscopy (XPS) (Fig. 1c), agree well with the earlier report.²⁶ The crystal morphology and elemental compositions of the GeS crystal, determined by scanning electron microscopy (SEM) and energy dispersive X-ray spectroscopy (EDS), manifest the layered nature of the GeS crystal (Fig. 1d) and the stoichiometric ratio of Ge : S $\sim 1 : 1$ (see ESI Fig. S1†).

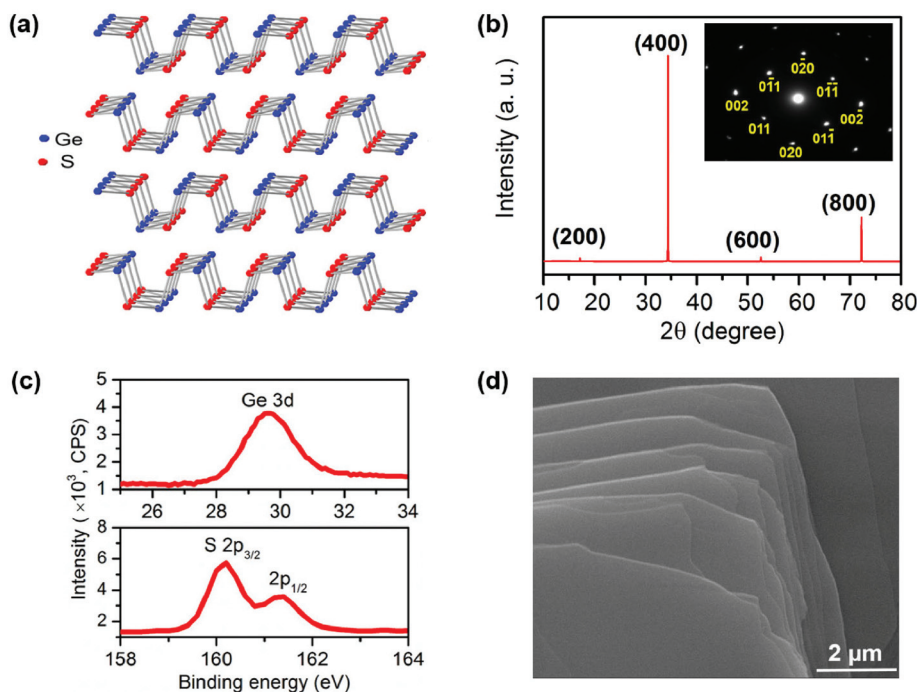


Fig. 1 (a) Crystal structure of GeS. (b) XRD spectrum of the as-synthesized bulk GeS crystal. The inset shows the SAED pattern of GeS nanosheets. (c) Observed XPS spectra of the bulk GeS crystal indicate the binding energies of Ge (3d) at 29.6 eV, S (2p_{3/2}) at 160.2 eV, and S (2p_{1/2}) at 161.4 eV. (d) SEM image of layered GeS nanosheets.

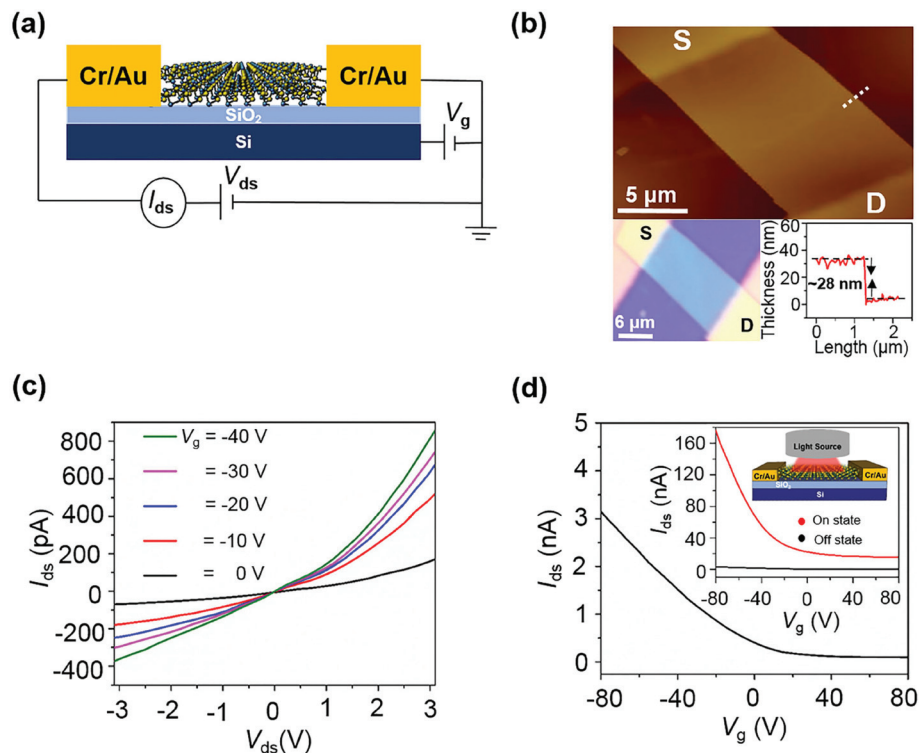


Fig. 2 (a) Illustration of a multi-layered GeS-FET. (b) AFM surface topograph (top panel) and optical image (bottom left panel) of the as-fabricated GeS-FET device. S and D represent the source and drain electrodes, respectively. The thickness of the GeS nanosheets, as scanned along the white dashed line in the top panel, is ~ 28 nm (bottom right panel). (c) The I_{ds} - V_{ds} characteristic curves of the multi-layered (~ 28 nm thick) GeS-FET measured at different V_g from 0 to -40 V. (d) The I_{ds} - V_g curve with V_g scanned from 80 to -80 V at $V_{ds} = 10$ V. The inset shows the transfer curves in the white-light illumination and dark states, in which the upper illustration represents the measurement of a multi-layered GeS-FET under illumination.

Multi-layered GeS-FETs, as illustrated in Fig. 2a, were fabricated following a standard lithographic procedure, where GeS nanosheets on a Si wafer of a 300-nm-thick SiO_2 dielectric layer were electrically connected to Cr/Au (5/70 nm) electrodes. Fig. 2b presents a surface topograph (top panel) and an optical image (bottom left panel) of the as-fabricated multi-layered GeS-FET with the thickness of the GeS nanosheets determined by atomic force microscopy (AFM) to be ~ 28 nm (bottom right panel), comprising ~ 50 successive GeS layers. Multi-layered GeS-FET devices fabricated with a semiconductive channel of ~ 28 nm thick GeS nanosheets were used throughout this study. The source-drain currents *vs.* source-drain voltage (I_{ds} - V_{ds}) of a multi-layered GeS-FET (Fig. 2c) were measured at different gate voltages ($V_g = -40$ – 0 V) in a probe station at 10^{-3} Torr and room temperature. The source-drain current *vs.* back-gate voltage (I_{ds} - V_g) curve (Fig. 2d) measured at $V_{ds} = 10$ V exhibits the p-type behavior of the GeS-FET because of the presence of Ge vacancies in the prepared GeS crystal.²⁷ The mobility of charge carriers in the GeS-FET can be deduced from $\mu = [L/(WCV_{ds})] \cdot (dI_{ds}/dV_g)$, where L and W are the length and width of the semiconductor channel of the GeS nanosheets, respectively, and C is the capacitance between the channel and the back gate per unit area ($C = \epsilon_0 \epsilon_r / d$; ϵ_0 is the vacuum permittivity, ϵ_r is the relative permittivity, and d is the

thickness of the SiO_2 dielectric layer). The field-effect mobility of the device is calculated to be $\sim 1.5 \times 10^{-3} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$; furthermore, the GeS-FET exhibits a high on/off ratio of $\sim 10^5$ (see ESI Fig. S2†). The presence of trap/defect states at the interface between the semiconductor channel and the SiO_2 dielectric layer of an FET device plays a dominant role in the charge carrier mobility.^{28,29} Reducing the number of such trap/defect states at the interface could substantially improve the device performances.^{30,31}

Next, we investigated the photocurrent generation in the GeS-FET by recording the I_{ds} - V_g curves under white-light illumination (shown in the inset of Fig. 2d). The I_{ds} - V_g measurements shown in the inset of Fig. 2d compare the photoresponses of the GeS-FET under illumination and in the dark. Upon illumination, the channel current was observed to increase dramatically relative to the dark state. Because of the excellent photoresponse of the multi-layered GeS-FET to white light, we conducted further photoelectrical studies of this device with various optical intensities. In this work, we carried out all of the photocurrent measurements under monochromatic illumination at 633 nm unless otherwise specified. Fig. 3a presents the I_{ds} - V_{ds} plots of a GeS-FET in the dark and under illumination with different intensities of 1.5, 3.5, 4.5, 15, 35, 45, and $153 \mu\text{W cm}^{-2}$, in which the photo-induced

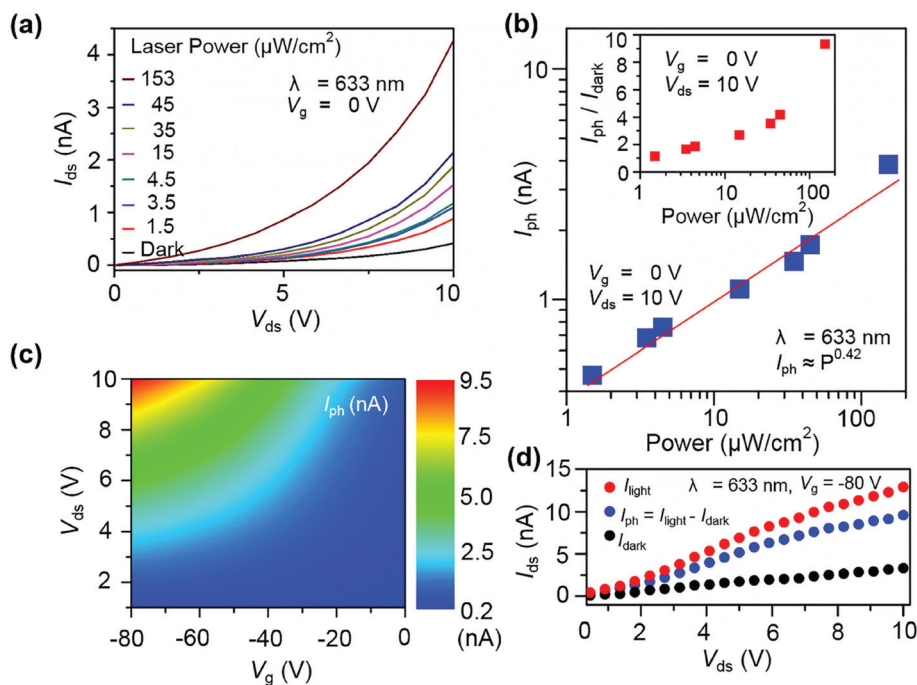


Fig. 3 (a) The $I_{\text{ds}}-V_{\text{ds}}$ curves of a multi-layered GeS-FET were investigated at $V_{\text{g}} = 0$ V with 633 nm excitation at laser powers of 1.5, 3.5, 4.5, 15, 35, 45, and 153 $\mu\text{W cm}^{-2}$. (b) The photocurrent (I_{ph}) is plotted as a function of the incident laser power (P) measured at $V_{\text{ds}} = 10$ V and $V_{\text{g}} = 0$ V; the data points (blue squares) were collected from (a). A linear fitting (the red solid line) to the experimental data (blue squares) reveals the dependence of the photocurrent on the incident laser power as $I_{\text{ph}} \approx P^{0.42}$. The inset shows the photosensitivity ($I_{\text{ph}}/I_{\text{dark}}$) as a function of incident laser power measured at $V_{\text{g}} = 0$ V and $V_{\text{ds}} = 10$ V; the data points (red squares) were collected from (a). (c) Mapping of photocurrent ($I_{\text{ph}} = I_{\text{light}} - I_{\text{dark}}$) as a function of V_{g} (from 0 to -80 V) and V_{ds} (from 0 to 10 V) under illumination of 1.5 $\mu\text{W cm}^{-2}$ at $\lambda = 633$ nm. (d) The $I_{\text{ds}}-V_{\text{ds}}$ curves of the multi-layered GeS photodetector measured in the illumination and dark states at $V_{\text{g}} = -80$ V, where the data points were taken from (c).

current gradually increases as the laser power increases. In Fig. 3b, we plot the generated photocurrent (I_{ph}) as a function of the incident laser power (P), where $I_{\text{ph}} = I_{\text{light}} - I_{\text{dark}}$ was calculated by subtracting the I_{ds} obtained in the dark (represented by I_{dark}) from that under illumination (denoted by I_{light}) at $V_{\text{ds}} = 10$ V. The photocurrent increases sublinearly following a power law of $I_{\text{ph}} \approx P^{0.42}$, which could be caused by the defects present in GeS and/or the adsorbed molecules/impurities at the GeS-SiO₂/Si interface.²⁹ The photosensitivity of a device is defined as the ratio of photo-generated current to dark current (*i.e.*, $I_{\text{ph}}/I_{\text{dark}}$). A detector of higher photosensitivity exhibits its ability of better noise rejection, where the photocurrent is clearly distinct from the dark current.³² In the inset of Fig. 3b, we present the increasing photosensitivity ($I_{\text{ph}}/I_{\text{dark}}$) as a function of the incident power to manifest the predominant sensitivity of a GeS photodetector to illumination. To gain further insight into the performance of the multi-layered GeS photodetector, the photocurrent ($I_{\text{ph}} = I_{\text{light}} - I_{\text{dark}}$) profile was measured under illumination at different V_{ds} (0–10 V) and V_{g} (-80 –0 V). The mapping in Fig. 3c shows an increase of I_{ph} by raising the negative V_{g} value with the concomitant rise of V_{ds} . Examining I_{ph} as a function of V_{ds} at a specific $V_{\text{g}} = -80$ V (Fig. 3d), the multi-layered GeS-FET possesses a low dark current, and the generated I_{ph} is dominated by the photo-generated charge carriers (*i.e.*, photocarriers)

upon illumination, making the multi-layered GeS photodetector as a good candidate for optoelectronic applications.

Photoresponsivity (R_{λ}) is an important parameter to judge the sensitivity of a photodetector under light stimulation. R_{λ} , which is defined as the photocurrent generated per unit power of the incident light (with wavelength λ) on an active illuminated area, is expressed as $R_{\lambda} = \Delta I_{\lambda}/(P_{\lambda}S) = [e\alpha/(h\nu)](\tau_{\text{p}}/\tau_{\text{t}})$, where ΔI_{λ} is the generated photocurrent, P_{λ} is the incident light power, S is the illuminated area, e is the elementary charge, t is the thickness of the channel, α is the absorption coefficient of the channel material at the photon energy of $h\nu$, and $\tau_{\text{p}}/\tau_{\text{t}}$ is the ratio of the life time of minority carriers (τ_{p}) to the transit time of majority carriers (τ_{t}).^{13,33} This relationship indicates that both α of the channel material and $\tau_{\text{p}}/\tau_{\text{t}}$ of the charge carriers are important factors that are responsible for the high photoresponsivity of a photodetector. Fig. 4a shows a 3D perspective of the measured R_{λ} (0.01–206 A W^{-1}) of a GeS-FET as a function of V_{ds} (0–10 V, $V_{\text{g}} = 0$ V) and P (153–1.5 $\mu\text{W cm}^{-2}$ at 633 nm). Fig. 4b plots R_{λ} in response to different incident laser power (P) values at 633 nm with a power law of $R_{\lambda} \approx P^{-0.6}$. Remarkably, the maximum R_{λ} of our multi-layered GeS photodetector reaches 206 A W^{-1} at 1.5 $\mu\text{W cm}^{-2}$ ($V_{\text{ds}} = 10$ V and $V_{\text{g}} = 0$ V). The high R_{λ} could be accounted for by the strong absorption nature of GeS in a wide spectral range^{19–22} together with the fast, effective separation of photo-

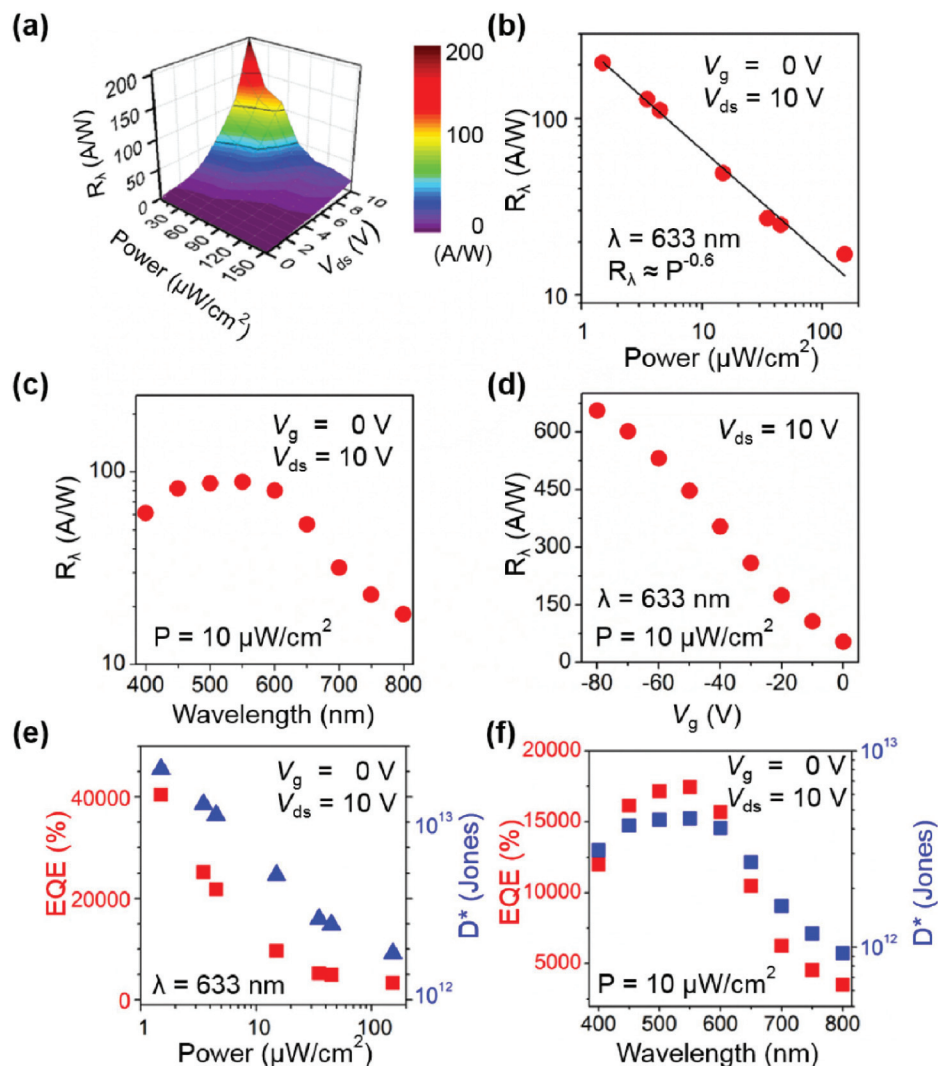


Fig. 4 (a) 3D photoresponsivity map of a multi-layered GeS photodetector. (b) Photo responsivity of the multi-layered GeS photodetector under various laser intensities at 633 nm. The power law of $R_{\lambda} \approx P^{-0.6}$ was determined from fitting the measured data. The GeS photodetector shows a maximal R_{λ} of 206 A W^{-1} excited at $1.5 \mu\text{W cm}^{-2}$, $V_g = 0 \text{ V}$, and $V_{ds} = 10 \text{ V}$. (c) The measured R_{λ} of the multi-layered GeS-FET device with respect to excitation wavelength. (d) R_{λ} of the multi-layered GeS photodetector as a function of V_g measured at $V_{ds} = 10 \text{ V}$ and $P = 10 \mu\text{W cm}^{-2}$. (e) The examined laser power dependences of EQE and D^* for the multi-layered GeS-FET at $V_g = 0 \text{ V}$ and $V_{ds} = 10 \text{ V}$. (f) The EQE and D^* of the multi-layered GeS photodetector were investigated as a function of excitation wavelength at $P = 10 \mu\text{W cm}^{-2}$, $V_g = 0 \text{ V}$, and $V_{ds} = 10 \text{ V}$.

carriers in the spatial 2D crystals of GeS nanosheets to defer electron-hole recombination and also by the large τ_1/τ_t , leading to a highly efficient photoresponsive conversion. In this study, we estimated $\tau_1/\tau_t \sim 1.4 \times 10^4$ for our GeS-FET device by calculating the τ_1/τ_t value from the above formula of R_{λ} with the assistance of computing $e\tau\alpha/(h\nu)$. Likewise, we also estimated the τ_1/τ_t values of multi-layered SnS_2 , multi-layered GeSe-, and GeS nanoribbon-based FET devices by extracting the available R_{λ} and $e\tau\alpha/(h\nu)$ values from reported literature data. These estimated τ_1/τ_t values, together with the channel thicknesses and lengths, of multi-layered GeS-, SnS_2 - and GeSe-FET devices and the GeS nanoribbon-based FET are listed in Table S1† for comparison. The GeS-FET performs substantially better in photoresponsivity as compared with GeS

nanoribbon-based (139.9 A W^{-1})³⁴ and the same group IV-VI family of GeSe- and SnS_2 -based photodetectors (see ESI Table S1†). The higher photoresponsivity of GeS nanosheet-based FETs, as compared to 1D GeS nanowire-based FETs, was also reported earlier²² and can be ascribed to the large surface area of the nanosheets to facilitate the more efficient generation of photo-excited carriers, resulting in higher photocurrent. In order to provide an insight into spectral response of the multi-layered GeS photodetector, we also measured the excitation wavelength dependences of R_{λ} at 400–800 nm, $P = 10 \mu\text{W cm}^{-2}$, $V_{ds} = 10 \text{ V}$, and $V_g = 0 \text{ V}$ (Fig. 4c). The R_{λ} value reaches a maximum of 89 A W^{-1} at $\lambda \sim 550 \text{ nm}$ and is generally in line with the trends of the observed absorption spectrum (see ESI Fig. S3†) and photoconductivity measurement³⁵ of a

bulk GeS crystal. Note that R_λ can be further increased by simply tuning the gate voltage (V_g) as shown in Fig. 4d,¹³ where the V_g -dependent R_λ was measured at $P = 10 \mu\text{W cm}^{-2}$, $V_{ds} = 10 \text{ V}$, and V_g (-80 – 0 V). Remarkably, the R_λ increased significantly from 54 to 655 A W^{-1} when V_g was swept from 0 to -80 V .

Next, we determined two other key parameters in photo-detection, EQE and D^* , for the multi-layered GeS photodetector. EQE is the measure of the photoconductive gain, which is the number of electrons detected per incident photon (*i.e.*, the photon-to-electron conversion), and can be expressed as $\text{EQE} = R_\lambda [hc/(e\lambda)]$, where R_λ is the photoresponsivity, h is Planck's constant, c is the speed of light, e is the elementary charge, and λ is the incident light wavelength. Based on the relationship of EQE to R_λ , the multi-layered GeS-FET possesses a high EQE of $\sim 4.0 \times 10^4\%$, which could be due to both high R_λ and the large surface-to-volume ratio of the 2D layered GeS nanosheets explained as follows. The processes for the photon-to-electron conversion start from the generation of free electron-hole pairs upon photon absorption. The spatial separation of photocarriers in the broad 2D crystal planes of GeS could preclude the electron-hole recombination. The aforementioned trap/defect states, which could be exacerbated further because of the high surface-to-volume ratio of the 2D GeS, were able to capture one type of photo-generated charge carrier while the other type escaped.^{36,37} As a result, the escaping photocarrier with prolonged lifetime can make several effective transits in the GeS channel between the electrodes, thus enhancing the photoconductive gain.^{38,39}

The D^* value, used to quantify the ability of a photodetector to detect weak optical signals, is defined by $D^* = (S \Delta f)^{1/2} / \text{NEP}$, where S and Δf are the effective area and electrical bandwidth of the photodetector, respectively, and NEP represents noise equivalent power. The NEP is an indicator of the minimum impinging optical power that a photodetector can distinguish a signal from noise. But, if the photodetector has a low value of NEP, then the above expression can be simplified as $D^* = R_\lambda S^{1/2} / (2eI_{\text{dark}})^{1/2}$, where R_λ , S , e , and I_{dark} are the photoresponsivity, effective area, elementary charge, and dark current, respectively.^{11,23,40–42} The high $D^* \sim 2.35 \times 10^{13}$ Jones (at $V_g = 0$ and $V_{ds} = 10 \text{ V}$) obtained for the multi-layered GeS photodetector is due to its low I_{dark} and high R_λ . The obtained D^* is comparable to the recent 2D material-based photodetectors, *e.g.*, GaS ($\sim 10^{13}$ Jones)¹¹ and In_2Se_3 ($\sim 10^{12}$ Jones)⁴² in the visible spectral regions and is close to the advanced commercial Si (10^{13} Jones)- and InGaAs (10^{12-13} Jones)-based photodiodes.²³ Fig. 4e presents the EQE and D^* of a multi-layered GeS photodetector measured with respect to the incident laser power. The excitation wavelength-dependences of EQE and D^* are also presented in Fig. 4f and the D^* vs. V_g measurement is shown in Fig. S4 of the ESI.†

Finally, time-resolved measurements were conducted to show the photoresponse and photoswitching stability of the multi-layered GeS photodetector in response to light illumination. The I_{ds} - t plot was recorded by exciting the multi-layered GeS device with a train of on-off illumination at 633 nm.

Fig. 5a shows the GeS-FET device in response to a single on-off cycle of illumination ($P = 12.7 \text{ mW cm}^{-2}$, $V_{ds} = 10 \text{ V}$, and $V_g = 0 \text{ V}$), which exhibits a fast rise of current (I_{ds}) under illumination and a sudden drop followed by a much slower relaxation after switching off the laser. While the rising time of the photo-generated signal was determined $7 \pm 2 \text{ ms}$ from a systematic measurement with several GeS-FET devices, the falling time consists of two components with a fast decay of $9 \pm 1 \text{ ms}$

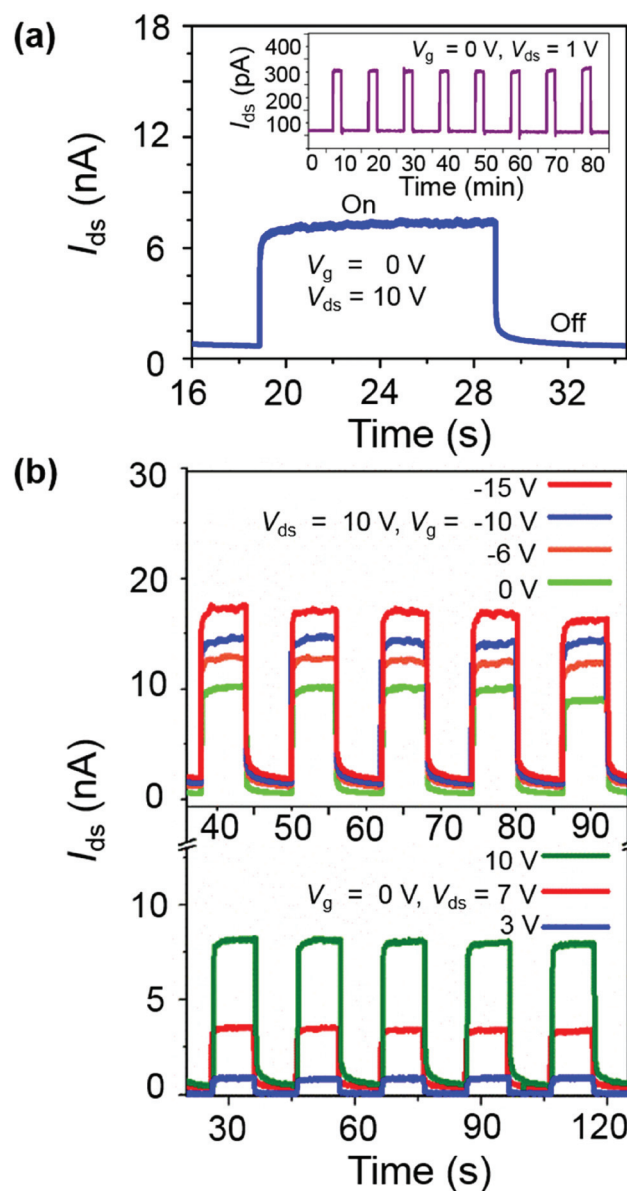


Fig. 5 (a) Time-resolved photoresponse of the multi-layered GeS photodetector was measured in response to a single on-off cycle of illumination at $P = 12.7 \text{ mW cm}^{-2}$ ($\lambda = 633 \text{ nm}$), $V_g = 0 \text{ V}$, and $V_{ds} = 10 \text{ V}$. The inset shows the photoswitching stability of a GeS photodetector in response to a train of pulsed illumination at $P = 12.7 \text{ mW cm}^{-2}$ ($\lambda = 633 \text{ nm}$), $V_g = 0 \text{ V}$, and $V_{ds} = 1 \text{ V}$. (b) Time-resolved photoresponse of the GeS photodetector was examined at various V_g ($= 0, -6, -10,$ and -15 V) and V_{ds} ($= 3, 7,$ and 10 V), respectively.

(corresponding to the earlier 50% decrease) and a slow relaxation of ~ 2 s (in the latter half of the decline). The response time of the multi-layered GeS photodetector is found to be faster than the recently reported GeS nanoribbon-based photodetector³⁴ by more than two orders of magnitude and is comparable to those of GaSe¹² and In₂Se₃-based photodetectors.⁴² The relatively long photocarriers lifetime, compared with traditional metal-semiconductor-metal photodetectors,⁴³ could again be due to the presence of trap/defect states in the GeS photodevice and/or on the surface of the SiO₂/Si substrate (wafer). Such trap/defect states could interrupt the electron-hole recombination, resulting in prolonging the lifetime of photocarriers.¹³ In response to consecutive photoswitching, the robustness and stability of our multi-layered GeS photodevice were examined by a train of pulsed illumination over a long period (~ 90 min) and a long on-off sequence (~ 100 cycles), as demonstrated by the $I_{\text{ds}}-t$ measurements presented in the inset of Fig. 5a and S5 of the ESI,[†] respectively. The time-resolved photoresponses of the multi-layered GeS photodevice were further tested at different $V_{\text{g}} = 0, -6, -10,$ and -15 V and $V_{\text{ds}} = 3, 7,$ and 10 V (shown in Fig. 5b), where the I_{ds} increases with the increasing amplitude of V_{g} or V_{ds} . These sustainable rise times and reproducible on-off shapes clearly demonstrate the device robustness, reproducible response, and switching stability of the multi-layered GeS photodetector.

Conclusion

In summary, single-crystalline GeS crystals were synthesized *via* a CVT method. High-quality multi-layered GeS nanosheets were obtained from a bulk GeS crystal *via* adhesive tape exfoliation. Multi-layered GeS-FET devices were fabricated following a standard lithographic procedure. The optoelectronic properties of the multi-layered GeS-FETs at room temperature were explored to render a very high $R_{\lambda} \sim 206$ A W⁻¹ (at $\lambda = 633$ nm, $P = 1.5$ $\mu\text{W cm}^{-2}$, $V_{\text{g}} = 0$ V, and $V_{\text{ds}} = 10$ V) and EQE $\sim 4.0 \times 10^4\%$, both of which are excellent among the recently reported 2D crystal-based photodetectors. Moreover, the value of $D^* \sim 2.35 \times 10^{13}$ Jones of the multi-layered GeS-FETs (measured at $V_{\text{g}} = 0$ V and $V_{\text{ds}} = 10$ V) is comparable to those of the advanced commercial Si- and InGaAs-based photodiodes. Finally, the multi-layered GeS photodetectors exhibit device robustness, photoswitching stability, and long-term durability. These extraordinary photoresponses allow the multi-layered GeS photodetectors to occupy a significant place in future optoelectronic applications.

Experimental section

Crystal growth

GeS crystals were grown from the CVT reaction in a three-zone horizontal furnace with the experimental conditions listed in Table 1. The furnace was constructed by using a quartz tube closed at one end with dimensions of 100 cm in length and

Table 1 The experimental conditions for the growth of GeS single crystals

Synthesized crystal	Reaction/growth temperature		Growth period	Dimensions of the as-grown crystal
	Reaction zone	Growth zone		
GeS	Zone 1 at 873 K Zone 2 at 850 K	Zone 3 at 800 K	100 h	$7.5 \times 6.5 \times 0.5$ mm ³

12 cm/10 cm in outer/inner diameter. For the crystal growth, high-quality quartz ampoules of 32 cm in length and 2 cm/1.8 cm in outer/inner diameter were filled with ~ 10 g samples containing the stoichiometric proportion (1 : 1) of Ge (99.999% purity) and S (99.999% purity) and were sealed at 10^{-5} Torr. After the sealed ampoules were placed in the three-zone horizontal furnace, the reaction and growth zones were slowly heated and later maintained at specific temperatures during the growth period. The experimental conditions for the growth of GeS single crystals are listed in Table 1.⁴⁴

Mechanical exfoliation

The as-synthesized single-crystal GeS flakes were exfoliated into multi-layered GeS nanosheets using a simple mechanical exfoliation method. Briefly, the bulk GeS crystal was placed on an adhesive tape, and the GeS layers were peeled off by slight rubbing and slicing. The rubbing-slicing process was repeated several times before transferring the GeS nanosheets to a Si wafer containing a 300 nm thick SiO₂ dielectric layer. The transferred GeS nanosheets on the Si wafer were inspected using an optical microscope (Olympus, BX 51M) equipped with a charge-coupled device (CCD) camera (Leica, DFC495). The thickness of the GeS nanosheets was estimated from their different interference colors in the optical image and was later determined from atomic force microscopy (AFM) (Veeco, BioScope SZ) scans.

Device fabrication

For the device fabrication of the multi-layered GeS-FETs, a transmission electron microscope (TEM) copper grid was mounted on the as-exfoliated GeS nanosheets on a Si wafer using a homemade micro-manipulator. The TEM grid acted as a shadow mask for the deposition of metallic electrodes (5 nm Cr/70 nm Au) in a thermal evaporator to fabricate the multi-layered GeS-FETs.

Optoelectronic characterization

The optoelectronic properties of the bulk GeS crystals, the as-exfoliated multi-layered GeS nanosheets, and the as-fabricated multi-layered GeS-FETs were characterized using various techniques. The lattice constant and crystal structure were examined with an X-ray diffractometer (X'Pert PRO-PANalytical, CuK α radiation). The morphology and crystallinity were investi-

gated by scanning electron microscopy (SEM) (FEI, Nova 200) and transmission electron microscopy (TEM) (JEOL, JEM-2100F). The electrical and optical parameters of the multi-layered GeS devices were measured at room temperature in a probe station (Lakeshore, TTPX) at 10^{-3} Torr equipped with a source meter (Keithley, 2636A) and an optical system, including a He-Ne laser (JDS Uniphase, Novette 1507), a power meter (Ophir, Nova II), an optical beam shutter (Thorlabs, SH1), a xenon lamp (Newport, 66921), and a monochromator (Acton, SpectraPro-500). The photoresponse time measurements were conducted *via* a Semiconductor Devices Parameter Analyzer (Agilent Technologies, B1500 A).

Author contribution statement

R. K. U. designed the experiments. R. K. U. and Y. Y. L. carried out electrical and optical measurements. C. J. K. performed spectral experiments with wavelength dependence. R. K. U. and S. R. T. fabricated field-effect transistor devices. R. S. and F. C. C. grew GeS crystals. K. M. B., A. A., K. Y., and R. J. M. assisted the experiments of electrical and optoelectronic characterization. C. R. L. conducted real-time photoresponsive measurements. R. K. U. wrote the manuscript. Y. T. C. is the corresponding author in charge of the project design, experimental data analysis, and manuscript writing.

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